

Part 1 : 1. References Electromagnetics theory [1] Cheng D. ' Fundamental Engineering Electromagnetics' 1st edn. Laser modes Examination of the laser output with a spectrometer of very high resolving power, such as the scanning Fabry–perot interferometer, reveals that it consists of a number of discrete frequency components (or very narrow spectral lines). The photoelectric effect, which is the emission of electrons from the surfaces of solids when irradiated, was explained by Einstein in 1905. He suggested that the energy of a light beam is not spread evenly, but is concentrated in certain regions, which propagate like particles. They are simplest form and easy to generate; arbitrary periodic time functions can be expanded into Fourier series of harmonic sinusoidal components; and transient nonperiodic functions can be expressed as Fourier integral. High order TEM modes or Hermite–Gaussian modes : We know the usual starting point for the derivation of laser beam propagation modes is solving the scalar Helmholtz equation (wave equation) within the paraxial approximation. However, while the wave theory, as we shall see below, provides an explanation of optical phenomena such as interference and diffraction, it fails completely when applied to situations where energy is exchanged, such as in the emission and absorption of light and the photoelectric effect. Field vectors that vary with space coordinates and are sinusoidal functions of time can be mathematically represented by vector phasors that depend on space coordinates and time as follows
$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}(\mathbf{r}) e^{i(kz - \omega t)}$$
 In contrast to this, the stimulated emission process results in coherent radiation since the waves associated with the stimulating and stimulated photons have identical frequencies, are in phase, have the same state of polarization and travel in the same direction. (22) And similarly for the other two components of \mathbf{E} and the three components of \mathbf{H} . The electric and magnetic fields vibrate perpendicularly to one another and perpendicularly to the direction of propagation as illustrated in Fig. 3 that is, light waves are transverse waves. discrete frequency components, they are not concerned with laser propagation where the longitudinal modes all contribute to a single 'spot' of light in the laser output, whereas the transverse modes discussed below may give rise to a pattern of spots in the output. Source of light It is well-known that when an electron in an atom undergoes transitions between energy states or levels it either emits or absorbs a photon, which can be described in terms of a wave of frequency ω . These are : (a) the spontaneous emission process in which the electron drops to the lower level in an entirely random way and (b) the stimulated emission process in which the electron is triggered to undergo the transition by the presence of photons of energy $2\hbar\omega$. Let us now examine the first phasor term on the right side of equation (42) and write
$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}(\mathbf{r}) e^{i(kz - \omega t)}$$
 For a cosine reference, the instantaneous expression for \mathbf{E} in equation (43) is
$$\mathbf{E}(\mathbf{r}, t) = \text{Re} \{ \mathbf{E}(\mathbf{r}) e^{i(kz - \omega t)} \}$$
 If such a source is located in an isotropic medium (such as free space) it will radiate uniformly in all directions, the wavefronts are thus a series of concentric spherical shells. B. Transverse modes Longitudinal modes are formed by the plane waves travelling axially along the laser cavity on a line joining the centers of the mirrors. The fundamental transverse electromagnetic mode (TEM₀₀) In the case of a fundamental transverse electromagnetic TEM₀₀ mode, the irradiance distribution across the beam is Gaussian, and so may write the electric field variation as
$$E_x = E_0 \exp(-x^2/w^2) \exp(-y^2/w^2) \exp(i(kz - \omega t))$$
 Square real part of the amplitude of Hermite–Gaussian modes, within the plane $z = 0$, for laser beams with beam waist radius $w_0 = 4 \times 10^{-3} \text{ m}$, by discharging an induction coil across a spark gap thereby

setting up oscillating electric and magnetic fields. Fig.3 Electromagnetic wave : the electric vector and magnetic vector vibrate in orthogonal planes and perpendicular to the direction of propagation. The remaining two equations (25) and (29) become

$$\frac{\partial^2 E_x}{\partial z^2} + k^2 E_x = 0 \quad (39)$$

$$\frac{\partial^2 H_y}{\partial z^2} + k^2 H_y = 0 \quad (40)$$

where partial derivatives have been replaced by ordinary derivatives since H_y and E_x are functions of only one variable, z . Thus H_y is the only nonzero component of corresponding H to the E in equation (43), and since $E_x = E_0 \cos(kz - \omega t)$, thus, the laser lineshape will have a finite wavelength (or frequency) spread i.e. they have a spectral width Δf . In general, there are many mathematical details related to this distribution and its properties, but what has been presented here is sufficient to clarify the spectral lineshape of laser light output.

Types of laser spectral lineshape broadenings

The spectral lineshape broadening is actually due to a number of external factors and internal atomic processes. Broadened laser transition line (or irradiance against frequency) (a) and (b) cavity modes (c) axial modes in the laser output. For any real laser cavity there will probably be waves travelling just off axis that are able to replicate themselves after covering a closed path such as Fig.9. These will also give rise to resonant modes, but because they have components of their electromagnetic fields which are transverse to the direction of propagation they are termed transverse electromagnetic (or TEM) modes. Visible light and Hertzian waves are part of the electromagnetic spectrum which, as we can see from Table 1, extends approximately over the wavelength range of 10^{-11} to 10^2 m. Planck's hypothesis did not require that the energy should be emitted in localized bundles and it could, with difficulty, be reconciled with the electromagnetic wave theory. For our purposes it is sufficient to accept that in many experiments, especially those involving the exchange of energy, the particle (photon or quantum) nature of light dominates the wave nature. There is nothing mystical in this, as the electron would undergo this process sooner or later spontaneously : the transition is simply initiated by the presence of stimulating photon.

Pure plane wave

A. Summary of Maxwell's equations : The results of combining Faraday's law, Ampere's law and Gauss' law are referred to as Maxwell's equations. (p4)

C. Time-Harmonic field Maxwell's equations and all the equations derived from them so far in this work hold for electromagnetic quantities with an arbitrary time dependence. (30)

D. Uniform plane wave A uniform plane wave is a particular solution of Maxwell's equations with E (and H) assuming the same direction, same magnitude, and same phase in infinite planes perpendicular to the direction of propagation. (45) It is clear that equation (45) represents a traveling wave and describes a perfectly monochromatic plane wave of infinite extent propagating in the positive z direction. (47) We can see, the second phasor term on the right side of equation (42) $E_0 \cos(kz + \omega t)$, represents a sinusoidal wave traveling in the $(-z)$ direction with the same velocity. This can be simply seen practically in both emission and absorption processes and if, for example, we were to measure the transmission (or emission) as a function of frequency for transition between the energy states E_1 and E_2 , we would obtain a probability distribution. Changing u corresponds to moving the curve to another position (translating it), and for $u = 0$ it is symmetric with respect to the ordinate (i.e. vertical direction), as shown in Fig.6.A.

Longitudinal modes The two mirrors of the laser form a resonant cavity and standing wave patterns are set up between the mirrors in exactly the same way that standing waves develop on the string. The modes of oscillation of the laser cavity will consist, therefore of a large number of frequencies, each given by

equation (73) and separated by $cL/2$, as shown in Fig.8. It should be appreciated, however, that while all the integers n give possible axial cavity modes only those which lie within the gain curve or laser transition line will actually oscillate. Fig. 10 shows the typical variation of w , with position, within a cavity formed by two concave mirrors of radius of curvature r_1 and r_2 separated by L .

Nature of light During the seventeenth century two emission theories on the Nature of light were developed, the wave theory of Hooke and Huygens and the corpuscular theory of Newton. Then in 1864 Maxwell combined the equation of electromagnetism in a general form and showed that they suggest the existence of transverse electromagnetic wave. Maxwell theory suggested the possibility of producing electromagnetic waves with a wide range of frequencies (or wavelengths). In 1887 Hertz succeeded in generating non-visible electromagnetic waves, with a wavelength of the order of μ m. As the absorption transition, in common with stimulated emission, can only occur in the presence of photon of appropriate energy, it is often referred to as stimulated absorption.

D E H J (11) B. Waves equations : First we derive the wave equation that governs the propagation of all electromagnetic wave. Each equation is composed of three scalar differential equations in terms of the components of the vectors. This simplified diagrams and mathematical descriptions but we should always remember that there is also a magnetic field component which behaves in similar way to the electric field component. We may arbitrarily assume the direction of E to be in the positive x direction; that is $\vec{E} = E_x \hat{x}$ (31) This x component of E is a function of only z since the field is to be uniform over the xy plane x and y is thus independent of x and y coordinates. $\nabla^2 E_x = -\mu_0 \epsilon_0 \frac{\partial^2 E_x}{\partial z^2}$ (41) $E_x = E_0 \exp(ikz) \exp(-i\omega t)$ (42) where E_0^+ and E_0^- are arbitrary constants that must be determined by boundary condition. Thus if we fix our attention on a particular point (a point of particular phase) on the wave, we set $\cos(kz - \omega t) = \cos(kz - \omega t + 2\pi n)$ or $\omega t - kz = \text{const} - 2\pi n$, from which we obtain $\frac{\partial \omega}{\partial k} = \frac{\omega}{k}$. Thus, for example have plane wave propagating in direction yz to the z axis with its wavefront normal to the yz plane, we can write $\vec{E} = E_0 \exp(ikz) \exp(-i\omega t) \hat{x}$ (57) $\vec{r} = r \hat{r} = r \sin\theta \hat{\theta} + r \cos\theta \hat{z}$ (58).

Laser line shapes In deriving the expression for the propagation of plane wave actually represents the ideal case. It implicitly assumes that all the atoms in either the upper or lower levels would be able to interact with the perfectly monochromatic wave with lineshape $\delta(\omega - \omega_0)$. Although, the spectral width of a laser output can be much less than that of ordinary light due to the spontaneous emission process, it cannot really be considered monochromatic wave. It is considered the most important continuous distribution because in applications many random variables are normal random variables, (that is, they have a normal distribution) or they are approximately normal or can be transformed into normal random variables in relatively simple fashion. Furthermore, the normal distribution is a useful approximation of more complicated distributions, and it also occurs in the proofs of various statistical tests. The normal distribution, also known in physics studies as Gaussian distribution, is defined as the distribution with density (or the probability density function) $f(x) = \frac{1}{\sigma\sqrt{2\pi}} \exp\left(-\frac{x^2}{2\sigma^2}\right)$. It is also a useful approximation of more complicated distributions, and it also occurs in the proofs of various statistical tests. The processes involved may be: (1) collision or (2) electromagnetic or (3) just the uncertainty broadening associated with the spontaneous lifetime. is the linewidth (full-width half maximum), that is the separation between the two points on the (frequency) curve where the function falls to half of its peak value which occurs at frequency $\omega_0 \pm \frac{\Delta\omega}{2}$. Practically, the inhomogeneous broadening

mechanisms lead to a Gaussian lineshape which may be written in terms of frequency as $\frac{1}{\Delta \nu} \exp\left(-\frac{2}{\Delta \nu} \left(\nu - \nu_0\right)^2\right)$ (74) As equation (74) is independent of ν_0 , the frequency separation of adjacent modes is the same irrespective of their actual frequencies. This is not accident but merely a direct consequence of the requirement that the mode be self replication as the light energy flows backwards and forwards between mirrors. Subsequent observations by Young, Malus, Euler and others lent support to the wave theory. On the other hand, for experiments involving interference and diffraction, where light interacts with light, the wave nature dominates. Let us consider the electron transitions which may occur between the two energy levels of the hypothetical atomic system shown in Fig 1. Under normal circumstances we do not observe the stimulated emission process because the probability of the the spontaneous emission is much greater than that of the stimulated emission. Because spontaneous radiation from any atom is emitted at random, the radiation emitted by a large number of atoms will clearly be incoherent. Energy level diagram illustrating (a) absorption, (b) spontaneous emission and (c) stimulated emission. Associated with Maxwell's equations, we have equation of continuity (or conservation of charge) $\nabla \cdot \mathbf{j} + \frac{\partial \rho}{\partial t} = -\nabla \cdot \mathbf{E}$ (16) We similarly obtain, by taking the curl of equation (13) and substituting equation (12), $\nabla \times \nabla \times \mathbf{u} = -\nabla^2 \mathbf{u} - \frac{1}{c^2} \frac{\partial^2 \mathbf{u}}{\partial t^2}$ In describing optical phenomena we often omit the magnetic field vector. (24) where $\mathbf{E}(\mathbf{r}, t) = \mathbf{E}_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t)$ and $\mathbf{H}(\mathbf{r}, t) = \mathbf{H}_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t)$ are the value of the electric and magnetic fields at the point \mathbf{r} at time t , E_0 and H_0 are the amplitudes of the electric and magnetic waves, (26) Expanding equations (25) and (26) in terms of components, the wave equations for the phasor components of the field vector become $\nabla^2 E_x = -k^2 E_x$, $\nabla^2 E_y = -k^2 E_y$, $\nabla^2 E_z = -k^2 E_z$ Strictly speaking a uniform plane wave does not exist in practice because a source infinite in extent would be required to create it, and practical wave sources are always finite in extent. Here we should mention that, equation (45) can also be expressed using a sine rather than a cosine function, or alternatively using complex exponentials. $\mathbf{u} = \mathbf{u}_0 \exp(i(\mathbf{k} \cdot \mathbf{r} - \omega t))$ (46) Equation (46) assures us that the the velocity of propagation of a equiphase front (the phase velocity) is equal to the velocity of light. As we know, it is impossible in practice to produce perfectly monochromatic waves, we often have the situation where a group of wave of closely similar wavelength is moving such that their resultant forms a packet. $\mathbf{y} = \int_{k_1}^{k_2} \cos(ky - \omega t) dk$ (59) By following the same analysis of one dimensional plane wave, hence we can write equation (56) in this case as $\frac{1}{\Delta \nu} \exp\left(-\frac{2}{\Delta \nu} \left(\nu - \nu_0\right)^2\right) \cos(\nu t - \mathbf{k} \cdot \mathbf{r})$ Intensity energy time area = is proportional to the square of the amplitude, there is an inverse-square-law decrease in irradiance. Likewise, it is a good check for robust techniques that are designed to work well under a wide variety of distributional assumptions. Practically, the homogeneous broadening mechanisms lead to a Lorentzian lineshape which may be written in terms of frequency as $\frac{1}{\Delta \nu} \frac{1}{1 + 4(\nu - \nu_0)^2 / \Delta \nu^2}$ In order to comply with what is required in laser light because the broadening in it occurs to the frequency (or wavelength). In some laser books, this lineshape is called Doppler frequency distribution because it is source of inhomogeneous broadening. In order to comply with what is required in laser light because the broadening in it occurs to the frequency (or wavelength). In this case electric field distributions are essentially given by the product of a Gaussian function and a Hermite polynomial, apart from the phase term as follows: $E(x, y, z, t) = u_0 \exp(i(\mathbf{k} \cdot \mathbf{r} - \omega t)) \exp\left(-\frac{x^2}{2\sigma^2} - \frac{y^2}{2\sigma^2}\right) H_n\left(\frac{x}{\sigma}\right) H_m\left(\frac{y}{\sigma}\right)$ and u_0 yielded a value for c in very close agreement with the value of the speed of light in vacuum measured independently. Maxwell therefore proposed that light was an electromagnetic wave having a speed of $c = 3 \times 10^8$ m/s, a frequency of some 14×10^{14} Hz and a wavelength of about 2×10^{-7} m. Planck

